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#### HALOGAS: the properties of extraplanar HI in disc galaxies

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#### **ABSTRACT**

We present a systematic study of the extraplanar gas (EPG) in a sample of 15 nearby late-type galaxies at intermediate inclinations using publicly available, deep interferometric HI data from the HALOGAS survey. For each system we mask the HI emission coming from the regularly rotating disc and use synthetic datacubes to model the leftover 'anomalous' HI flux. Our model consists of a smooth, axisymmetric thick component described by 3 structural and 4 kinematical parameters, which are fit to the data via a Bayesian MCMC approach. We find that extraplanar HI is nearly ubiquitous in disc galaxies, as we fail to detect it in only two of

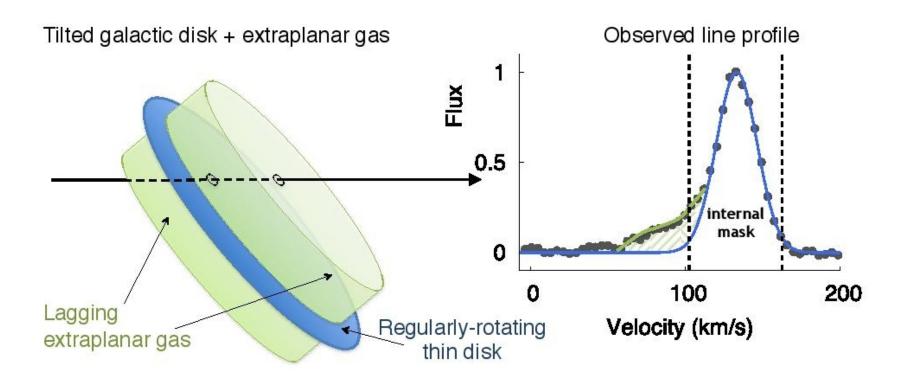
### Выборка из 15 спиральных галактик

**Table 1.** Physical properties of galaxies in our sample, from H11 and Heald et al. (2012). We also list the median kinematical inclination (INC<sub>BB</sub>) and position angle (PA<sub>BB</sub>) found with <sup>3D</sup>BAROLO and used in our EPG modelling.

UGC	Other name	Type	dist.	INC <sub>H11</sub>	$INC_{BB}$	$PA_{BB}$	D <sub>25</sub>	$M_B$	$v_{\rm rot}$	SFR
			(Mpc)	(°)	(°)	(°)	(')	(mag)	$(km s^{-1})$	$(M_{\odot}  yr^{-1})$
1256	NGC 0672	SBcd	7.6	70	67.6	64.2	6.4	-18.65	130.7	0.23
1913	NGC 0925	SABd	9.1	54	57.9	284.7	11.3	-19.66	102.4	0.77
1983	NGC 0949	SAd	11.3	52	52.5	160.8	3.5	-17.85	90.9	0.31
2137	NGC 1003	SAcd	11.6	67	70.4	276.3	6.3	-18.61	95.5	0.40
3918	NGC 2403	SAcd	3.2	62	62.5	124.6	23.8	-19.68	121.9	0.6
4284	NGC 2541	SAcd	12.0	67	63.8	171.8	7.2	-18.37	92.1	$0.55^{a}$
5572	NGC 3198	SBc	14.5	71	70.0	214.3	8.8	-19.62	148.2	1.1
7045	NGC 4062	SAc	16.9	68	67.1	100.1	4.5	-18.27	140.5	0.67
7353	NGC 4258	SABbc	7.6	71	74.0	331.9	17.1	-20.59	208.0	1.7
7377	NGC 4274	SBab	19.4	72	71.3	279.8	6.5	-19.22	239.9	1.2
7539	NGC 4414	SAc	17.8	50	53.9	159.7	4.5	-19.12	224.7	4.2
7591b	NGC 4448	SBab	9.7	71	73.5	94.4	3.8	-18.43	221.6	0.056
7766	NGC 4559	SABcd	7.9	69	68.0	323.1	11.3	-20.07	113.4	0.69
8334	NGC 5055	SAbc	8.5	55	65.2	99.4	13.0	-20.14	215.5	2.1
9179	NGC 5585	SABd	8.7	51	50.4	48.2	5.5	-17.96	79.1	0.41

<sup>&</sup>lt;sup>a</sup> based on TIR+UV measurements from Thilker et al. (2007) and re-scaled to the distance used in this work. Heald et al. (2012) report only an upper limit based on a non-detection with IRAS  $25\mu$ m.

### Из полного куба данных WSRT вычитается тонкий диск HI



**Fig. 1.** A sketch of an observation of extraplanar gas in a galaxy seen at intermediate inclination. The typical line profile (along the kinematic major axis of the galaxy) will be composed by a nearly Gaussian part coming from the thin disc with overlaid a tail at low rotation velocity produced by the lagging EPG layer. The width of the disc emission is roughly symmetrical, produced by gas turbulence and well fitted by a Gaussian function (blue solid line). The EPG separation is achieved by masking a portion of the profile with substantial contribution from this Gaussian function ('internal mask' region).

### К тому, что осталось, применяют модель толстого диска:

$$\Sigma(R) = \Sigma_0 \left( 1 + \frac{R}{R_g} \right)^{\gamma} \exp \left( -\frac{R}{R_g} \right),$$

$$\rho(z) \propto \frac{\sinh(|z|/h)}{\cosh^2(|z|/h)}$$

### $ho(z) \propto rac{\sinh(|z|/h)}{\cosh^2(|z|/h)}$ , 3 CTDYKTYDLB M LADAMETDA,N4-KNHONATNUOCKNX

The EPG kinematics is described by four parameters: the vertical gradient in the gas rotational speed  $(dv_{\phi}/dz)$ , the velocities in the radial and in the vertical directions ( $v_R$  and  $v_z$ ), and gas velocity dispersion  $\sigma$ . Thus the EPG is allowed to rotate with a different speed with respect to the material within the disc (or to not rotate at all, for  $dv_{\phi}/dz \ll 0$ ). It can globally accrete onto or escape from the galaxy, can move in/out and have a different velocity dispersion. These simple kinematical parameters allow us

### Индивидуальные результаты:

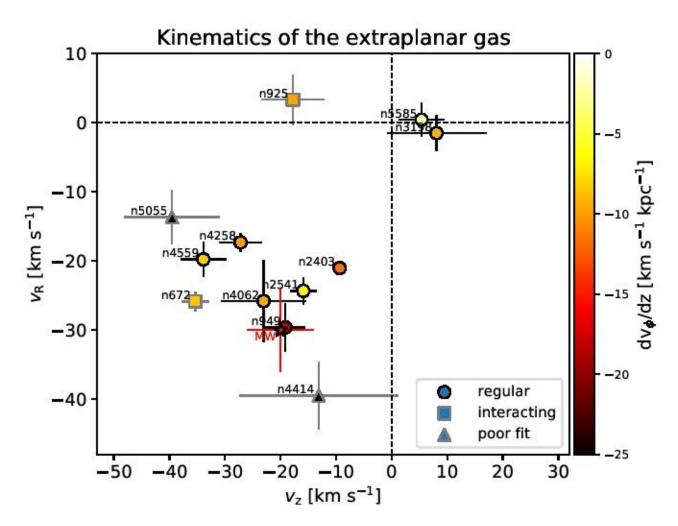
Table 4. Best-fit parameters and associated uncertainty for the EPG of the galaxies studied in this work.

Galaxy (1)	Beam (arcsec) (2)	Independent voxels (3)	$M_{\rm HI, EPG}$ (10 <sup>8</sup> ${ m M}_{\odot}$ ) (4)	$f_{\text{EPG}}$ (5)	R <sub>g</sub> (kpc) (6)	γ (7)	h (kpc) (8)	$ \frac{dv_{\phi}/dz}{(km s^{-1} kpc^{-1})} $ (9)	$v_R \ (\text{km s}^{-1}) \ (10)$	$(\text{km s}^{-1})$ (11)	$ \begin{array}{c} \sigma \\ (\text{km s}^{-1}) \\ (12) \end{array} $
NGC 0925 <sup>a</sup>	38x33	870	7.6	0.14	$0.9 \pm 0.1$	$13.3 \pm 2.1$	$2.0\pm0.1$	-9.7±0.7	$3.3 \pm 3.5$	$-17.8 \pm 5.6$	21.6±1.3
NGC 0949	39x35	190	1.6	0.27	$3.9 \pm 0.8$	$0.5 \pm 0.5$	$1.6\pm0.1$	$-18.0 \pm 1.4$	-29.6±3.5	$-19.1 \pm 3.5$	$27.7 \pm 1.6$
NGC 2403	30x29	4797	5.9	0.22	$2.3 \pm 0.1$	$2.6 \pm 0.3$	$0.9\pm0.1$	$-11.7 \pm 0.5$	-21±0.9	-9.4±1.2	$15.2 \pm 0.4$
NGC 2541	34x34	634	7.4	0.15	-	$4.9 \pm 1.0$	$2.8\pm0.1$	$-6.5\pm0.3$	$-24.3\pm1.9$	$-15.9\pm2.3$	$27.4 \pm 1.0$
NGC 3198	35x33	637	9.8	0.09	$2.1 \pm 0.5$	$9.8 \pm 2.8$	1.4±0.2	$-9.2 \pm 1.4$	$-1.5\pm2.6$	$8.1 \pm 8.9$	$23.0\pm1.2$
NGC 4062	39x34	148	2.1	0.11	-	$5.0\pm2.2$	$2.5\pm0.1$	$-9.4 \pm 1.0$	$-25.8 \pm 5.8$	$-23\pm7.5$	$34.6 \pm 2.0$
NGC 4258	33x33	1422	6.6	0.11	$6.6 \pm 1.1$	$2.3 \pm 0.4$	$1.4\pm0.1$	$-10.0\pm0.5$	$-17.3 \pm 1.3$	$-27.2\pm3.8$	24.8±0.7
NGC 4414 <sup>b</sup>	39x33	194	5.1	0.12	-	$-3.7\pm1.9$	$0.5\pm0.3$	$-58^{+24}_{-56}$	$-39.5 \pm 4.8$	-13.1±14.2	26.3±2.9
NGC 4559	41x32	926	6.0	0.13	$1.9 \pm 0.3$	$4.0 \pm 0.9$	$1.7 \pm 0.1$	-8.2±0.6	$-19.8\pm2.5$	$-33.9 \pm 4.1$	$21.3 \pm 1.2$
NGC 5055 <sup>b</sup>	36x33	491	2.6	0.04	$1.1 \pm 0.2$	$8.3 \pm 1.5$	$0.6\pm0.1$	$-40^{+8}_{-13}$	$-13.7 \pm 3.9$	$-39.5 \pm 8.6$	$20.2\pm2.4$
NGC 5585	34x33	368	3.0	0.12	$2.5^{+1.9}_{-0.9}$	$2.6 \pm 1.8$	$1.4\pm0.3$	-2.3±0.8	$0.4 \pm 2.4$	$5.4 \pm 4.0$	$18.5 \pm 1.0$

**Notes.** (1) NGC name; (2) FWHM of the synthesised beam; (3) Number of independent EPG voxels modelled, computed as the number of voxels in the masked dataset with intensity above twice the rms-noise divided by the number of voxels per resolution element (eq. 6); (4) H I mass of the EPG component, with an uncertainty of  $\sim 20\%$ ; (5) ratio between  $M_{\rm EPG}$  and the total H I mass; (6-7) surface density parameters (eq. 1), a '-' indicates that the parameter is unconstrained; (8) scale height (eq. 2); (9) vertical rotational gradient; (10) velocity in the direction perpendicular to the rotational axis (positive values mean outflow); (11) velocity in the direction perpendicular to the disc (positive values mean outflow); (12) velocity dispersion.

<sup>&</sup>lt;sup>b</sup> Poorly-fit systems, which are also warped. The fit is limited to (and  $f_{EPG}$  is computed within) the innermost 20 kpc.

## В основном, аккреция: радиальная и вертикальная



### Происхождение толстого диска HI: «Фонтанное»

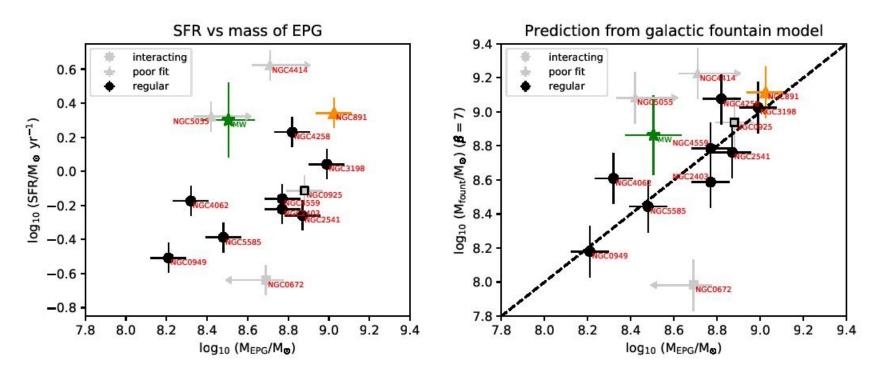


Fig. 11. Left panel: SFR vs EPG mass for our HALOGAS sample. Right panel: theoretical EPG mass predicted by a galactic fountain model (eq. 11 with  $\beta = 7$ ) vs that inferred from the data. The two panels span the same range ( $\sim 1.5$  dex) on both axes. Error-bars are determined by assuming a 10% uncertainty on the distances and a 20% uncertainty on the total SFR, in addition to the uncertainty on the  $R_{SFR}$  discussed in the text. Interacting galaxies and systems with a poorly fit EPG are shown with separate symbols. We have also included NGC 891( $M_{EPG}$  from Marinacci et al. 2010b) and the Milky Way, for which we used  $M_{EPG} = 3.2 \pm 1.0 \times 10^8$  M<sub> $\odot$ </sub> (MF11), and assumed fiducial values for SFR,  $R_{SFR}$  and  $v_{flat}$  of  $2 \pm 1$  M<sub> $\odot$ </sub> yr<sup>-1</sup>,  $2.5 \pm 0.5$  kpc and  $240 \pm 10$  km s<sup>-1</sup> respectively.

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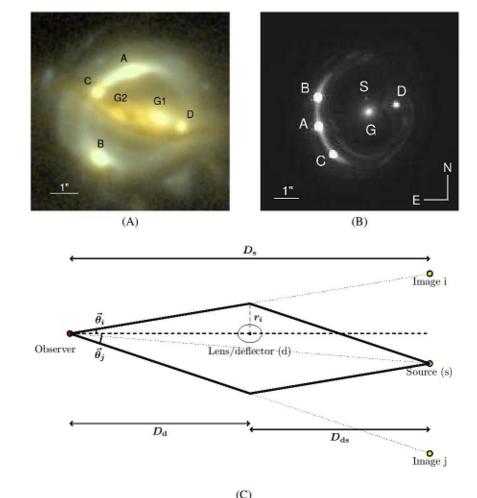
### A measurement of the Hubble constant from angular diameter distances to two gravitational lenses

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#### Резюме

The local expansion rate of the Universe is parametrized by the Hubble constant,  $H_0$ , the ratio between recession velocity and distance. Different techniques lead to inconsistent estimates of  $H_0$ . Observations of Type Ia supernovae (SNe) can be used to measure  $H_0$ , but this requires an external calibrator to convert relative distances to absolute ones. We use the angular diameter distance to strong gravitational lenses as a suitable calibrator, which is only weakly sensitive to cosmological assumptions. We determine the angular diameter distances to two gravitational lenses,  $810^{+160}_{-130}$  and  $1230^{+180}_{-150}$  Mpc, at redshifts of z=0.295 and 0.6304. Using these absolute distances to calibrate 740 previously-measured relative distances to SNe, we measure the Hubble constant to be  $H_0 = 82.4^{+8.4}_{-8.3} \,\mathrm{km}\,\mathrm{s}^{-1}\,\mathrm{Mpc}^{-1}$ .



# Включение в модель кинематики звезд в линзе очень улучшает точность

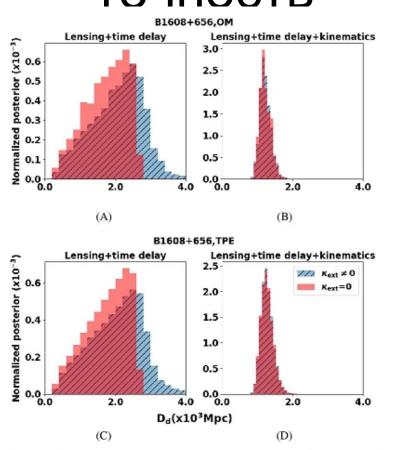


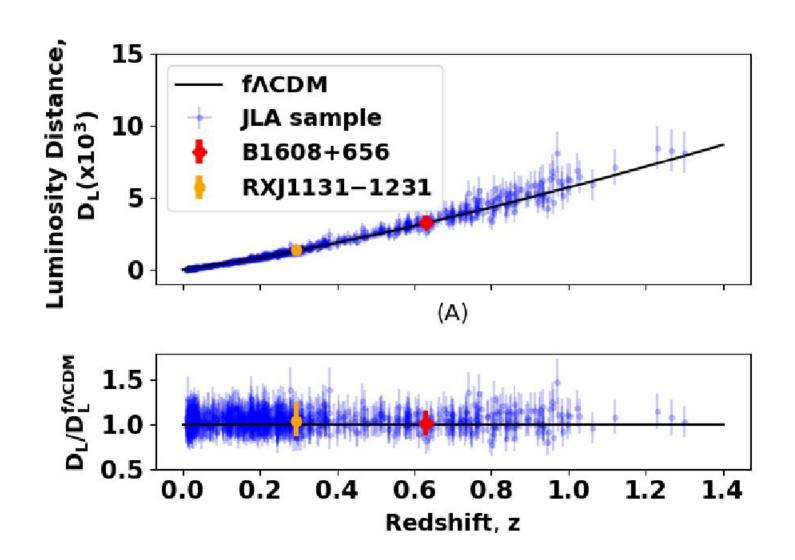
Figure 3: Normalized posterior probability distributions for the angular diameter distance to the lens B1608+656. (A, C) include lensing and time-delay information, while (B, D) include additionally the kinematics of the lens. The blue hatched distribution shows the results if the external convergence distribution is estimated by ray-tracing through the Millennium Simulation (36) (Fig. S4), while the red distribution is the result when the external convergence is set to zero. By including the kinematic information, the angular diameter distance becomes insensitive to  $\kappa_{\rm rest}$ .

### Измеренные расстояния до галактиклинз и их использование для калибровки SNI

Our analysis constrains the angular diameter distances to 12–20% precision per lens. We marginalize over the uncertainties in anisotropy models by merging two posterior probability distributions of OM and TPE models (Section S6). Our final measurements of the angular diameter distances are  $D_{\rm d}(z=0.6304)=(1.23^{+0.18}_{-0.15})\times 10^3$  Mpc for B1608+656, and  $D_{\rm d}(z=0.295)=(8.1^{+1.6}_{-1.2})\times 10^2$  Mpc for RXJ1131–1231.

We apply these distances as anchors to the 740 SNe in the Joint Light-curve Analysis (JLA, (37)) dataset, allowing us to constrain  $H_0$  and the SNe nuisance parameters (Section S7) simultaneously. We use the MontePython code (38) to perform a Markov Chain Monte Carlo analysis. Fig. 5 shows the resulting Hubble diagram, i.e. the absolute luminosity distances  $D_{\rm L} = (1+z)^2 D_{\rm d}$  as a function of redshifts for a flat  $\Lambda$ CDM model.

### Откалиброванная диаграмма Хаббла для сверхновых



### Полученная из нее постоянная Хаббла

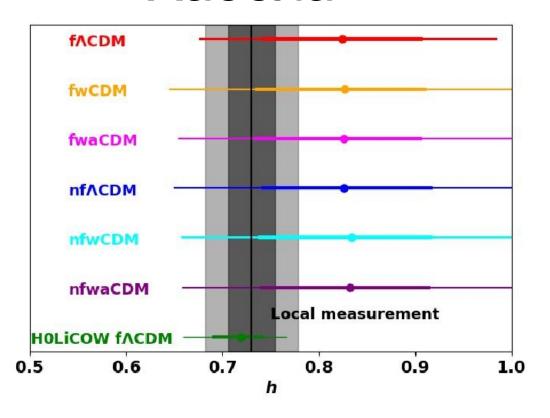


Figure 6: Constraints on the Hubble constant for six cosmological models. The gray shaded area is the constraint from the local distance ladder (4), while the green line is from three timedelay distances measured by the H0LiCOW collaboration (24). The thick and thin solid lines denote the 68% and 95% Confidence Levels of the joint fit to the SNe and the  $D_{\rm d}$  data. We emphasize that  $D_{\rm d}$  and  $D_{\Delta t}$  determined from the same lens are not correlated strongly because the uncertainty in the former is dominated by the kinematics and the latter by  $\kappa_{\rm ext}$  in the case of